

Magnetic domains tune faster out-of-equilibrium critical dynamics [1]

ESCUELA DE DOCTORADO DE MADRIE



I. González-Adalid Pemartín¹, E. Mompó², A. Lasanta^{3,4}, V. Martin-Mayor^{1,5}, and J. Salas^{2,4}

¹Departamento de Física Teórica, Universidad Complutense, Madrid, Spain

²Departamento de Matemáticas, Universidad Carlos III de Madrid, Leganés, Spain.

³Departamento de Álgebra, Facultad de Educación, Economía y Tecnología de Ceuta, Universidad de Granada, Ceuta, Spain.

⁴Grupo de Teorías de Campos y Física Estadística, Instituto Gregorio Millán, Universidad Carlos III de Madrid, Unidad Asociada al Instituto de Estructura de la Materia, CSIC, Spain ⁵Instituto de Biocomputación y Física de Sistemas Complejos (BIFI), Zaragoza, Spain

1. Introduction

Finding thermal equilibrium may be difficult, specially if one needs to do it fast. However, because reaching thermal equilibrium is a very important problem in industry, technology, and science, the non-equilibrium process are very studied, trying to answer the question:

How can we reach thermal equilibrium faster?

There are some examples of different non-equilibrium dynamics that allow a faster equilibration, like

• Mpemba effect [2],

Model
Ferromagnetic Ising model in 2D
$\mathscr{H} = -J \sum_{\langle \mathbf{x}, \mathbf{y} \rangle} \sigma_{\mathbf{x}} \sigma_{\mathbf{y}} , \qquad (1)$
with $\sigma_{\mathbf{x}} = \{\pm 1\}, \langle , \rangle$ the sum over nearest-neighbors,
a thermal bath at inverse temperature $\kappa = \frac{J}{K_{\rm B}T}$, energy
nits $J = 1$, and system size $L = 4096$. The system has a
nd phase transition at $\kappa_c\simeq 0.44068$, which a paramag-

3. Discussion

Real speed-up



- annealing techniques [3],
- stochastic reset.

In particular, Admit and Raz designed a strategy for system with time-scale separation [4]. Nevertheless, what happen with systems with a continuum time scale, e. g. a 2nd phase transition? So, in this work we try to answer the question:

What is the fastest way to reach thermal equilibrium in systems with a 2nd order phase transition?

netic phase (PM) for $\kappa < \kappa_c$, and a ferromagnetic phase (FM) for $\kappa_c < \kappa$. We know exactly $E_{eq} = E(t \rightarrow \infty)$ [5]. We use the Metropolis (MET) and Heat-Bath (HB) dynamics rules (*model*-A universality class [6]).

Observables

Finally, in order to study our system we consider:

the correlation function

 $C(\mathbf{r};t) = \frac{1}{I^2} \sum_{\mathbf{x}} \langle \sigma_{\mathbf{x}}(t) \sigma_{\mathbf{x}+\mathbf{r}}(t) \rangle,$

- the energy density $E(t) = -2C(\mathbf{r}_{\min}; t)$,
- and the coherence length $\xi(t)$.

Thermal protocols

Isothermal Protocol (1P)

- **1** Generate a disordered spin configuration.
- 2 Put the configuration in a thermal bath at
 - inverse-temperature κ (t = 0).
- **3** Evolve the configuration *t* Monte Carlo steps.





Figure 6: Energy density as a function of the total time elapsed since the beginning of 2P, t_{total} (MET). The width of the curves is twice the statistical errors (ξ_{start} increases from top to bottom). The horizontal lines correspond to E_{eq} multiplied by 1.001,1 and 0.999. **Inset:** coherence length during the 2P as a function of the time t.

Equilibration time



Figure 7: Equilibration time $t_{eq}^{0.1\%}$ (defined as the last time the system has an energy with a difference less than a 0.1% from E_{eq}) versus ξ_{start} for 2P. versus ξ_{start} for 2P. The width of the curves is twice the statistical errors. The discontinuity is due to the non-monotonic time behavior of $E(t_{total})$.



Figure 1: Cartoon of the system evolution using 1P. Schematic representation of the coherence length ξ .

Two-step Protocol (2P)

1 Generate a disordered spin configuration. 2 Put the configuration the FM phase. **3** Evolve the system until $\xi(t^*) = \xi_{\text{start}}$. • At t = 0 heat back to the PM phase.

Figure 2: Coherence length ξ as a function of time t for 1P and MET (the width of the curves is twice the statistical error; κ increases bottom to top). **Inset:** Comparison of the HB and MET for $\kappa < \kappa_c$.



2. Results

Isothermal evolution



Exponential speed-up



4. Conclusions

- Precooling may produce a faster equilibration at high T.
- The speed-up is driven by the magnetic domains in FM phase.
- An exponential speed-up is possible if the size of the magnetic domains is a well-defined fraction of the equilibrium correlation length in the high T phase.
- This will valid for the whole *model-A* universality class.
- We can use this mechanism with our *two-step* protocol.

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